

A Parametric Study of the SASI Comparing General Relativistic and Non-Relativistic Treatments

Samuel J. Dunham

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Paper

A Parametric Study of the SASI Comparing General Relativistic and Non-Relativistic Treatments*

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AND KELLY HOLLEY-BOCKELMANN ,^{1,5}

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Overview

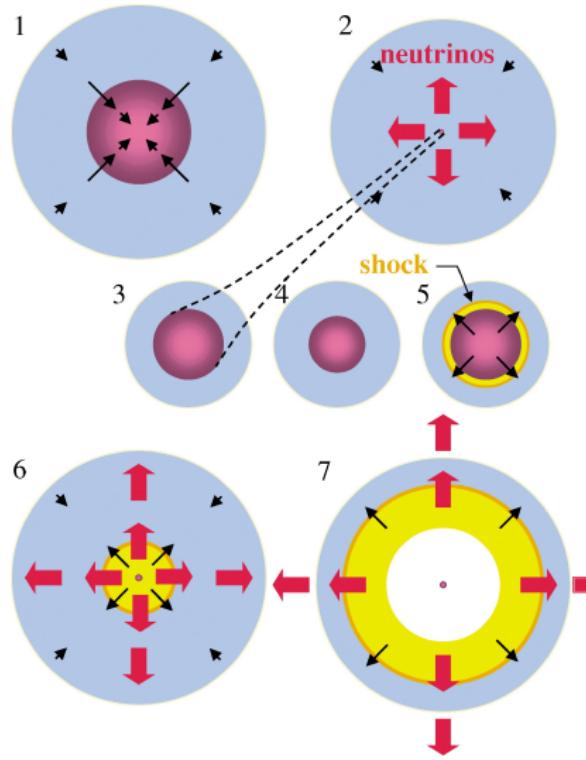
1 Background

2 Physical Model

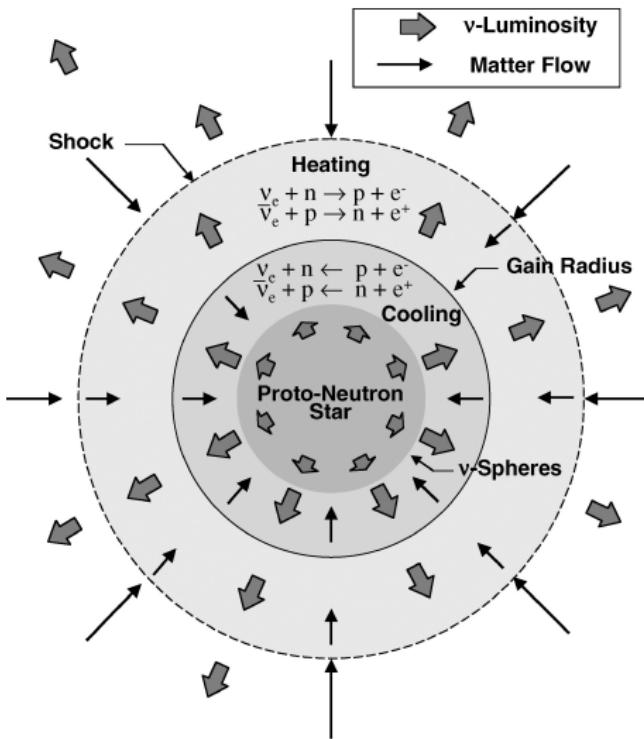
3 Steady-State Solutions

4 Results

5 Summary, etc.



Mezzacappa A. 2005.
Annu. Rev. Nucl. Part. Sci. 55:467–515



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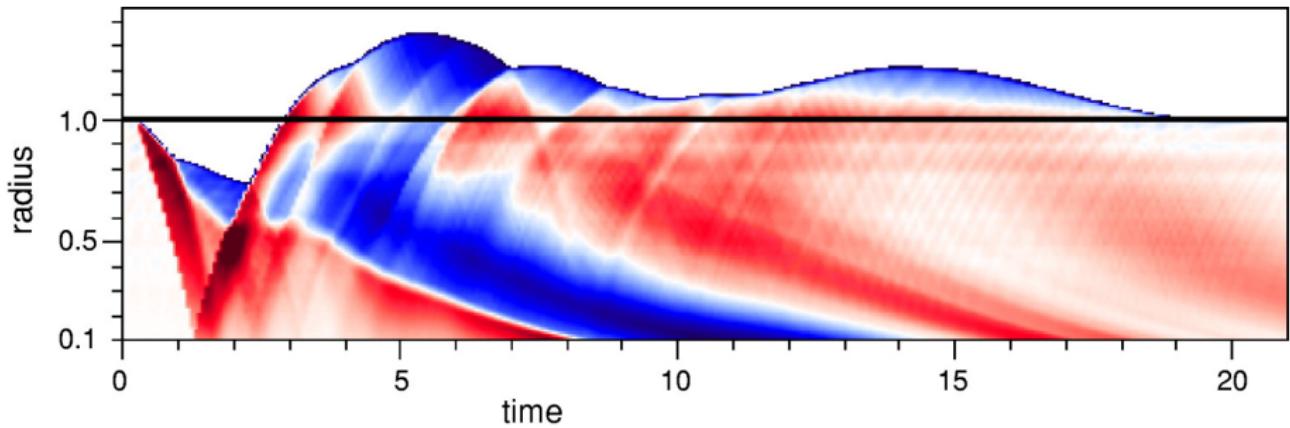


Figure: Equilibrium (white), under- (blue), and over- (red) pressure (Blondin et al., 2003).

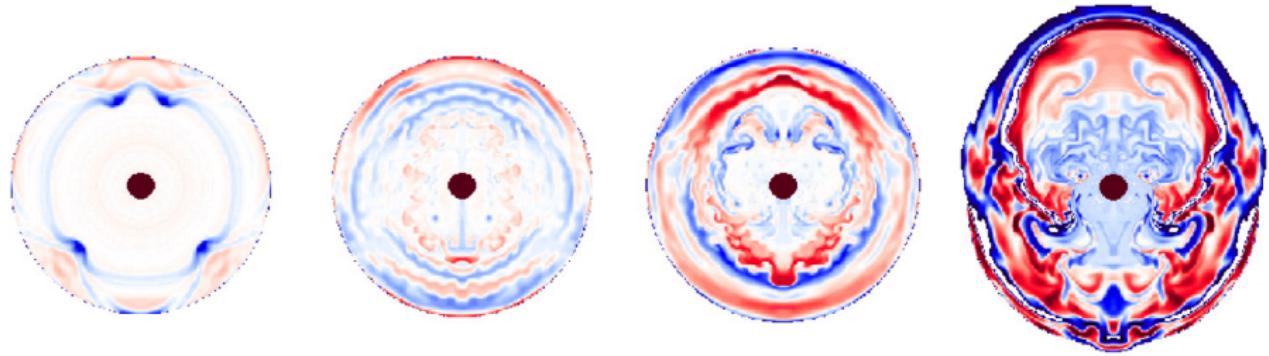
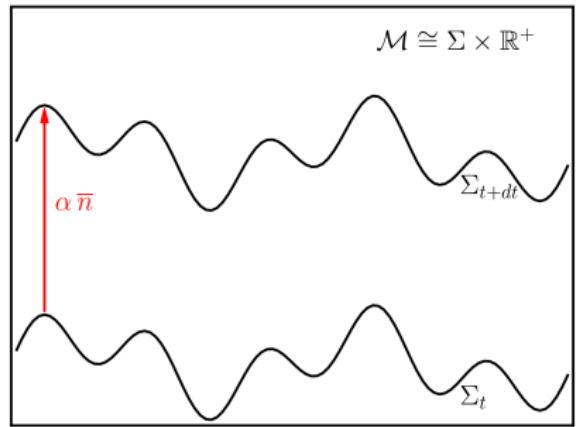


Figure: Equilibrium (white), under- (blue), and over- (red) entropy (Blondin et al., 2003).

$d+1$ Decomposition

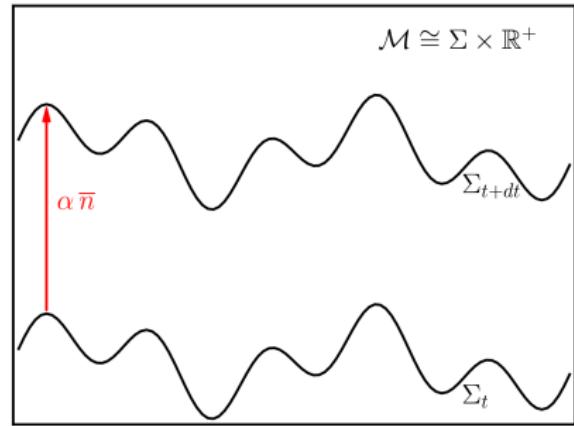
$$ds^2 = g_{\mu\nu} dx^\mu dx^\nu$$



$d+1$ Decomposition

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\underline{g} : spacetime metric on \mathcal{M}

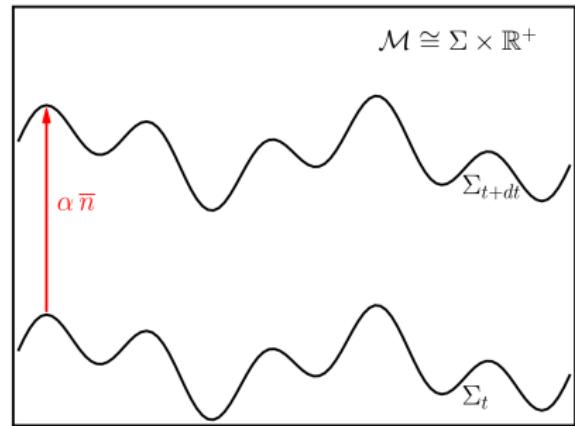


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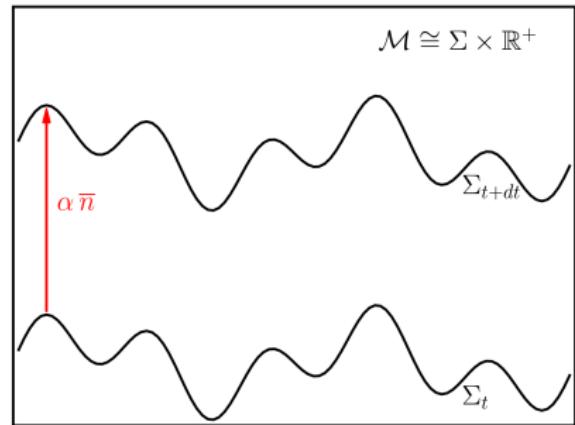
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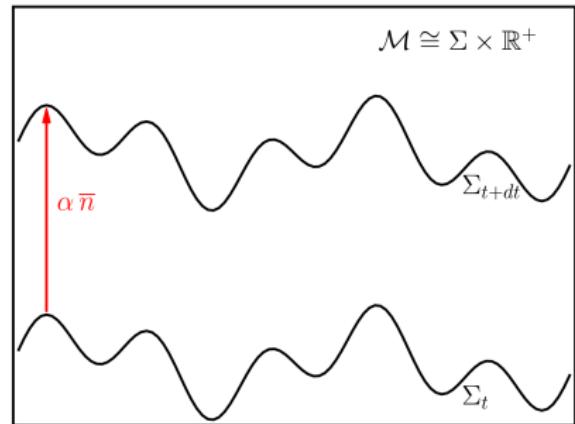
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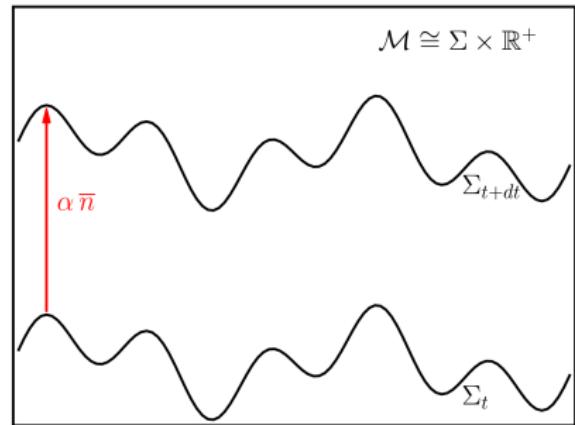
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$$\underline{\underline{(g(\bar{n}, \bar{n}) = \bar{n} \cdot \bar{n} = -1)}}$$



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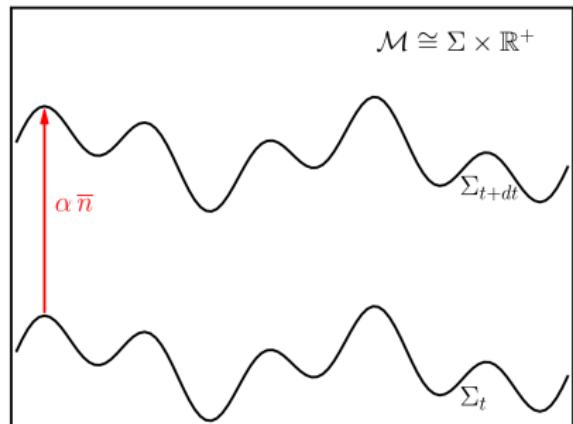
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$$(\underline{g}(\bar{n}, \bar{n}) = \bar{n} \cdot \bar{n} = -1)$$

(Also, \underline{K} : Extrinsic curvature)



Conformally-Flat Condition

Wilson et al. (1996); Cordero-Carrión et al. (2009)

$$\gamma_{ij} = \psi^4 \bar{\gamma}_{ij}$$

$1 \leq \psi < 2$: Conformal factor

$$\bar{\gamma}_{ij} = \text{diag}(1, r^2, r^2 \sin^2 \theta)$$

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(Also, $K := \text{Tr}_{\gamma_{ij}} (\underline{\underline{K}}) = \partial_t K = 0$)

Isotropic Coordinates ($c = G = 1$)

Baumgarte and Shapiro (2010)

$$\alpha(r) = \frac{1 - r_{\text{Sc}}/r}{1 + r_{\text{Sc}}/r}$$

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$$\beta^i = 0$$

$$K_{ij} = 0$$

Fluid Equations

Units defined such that $c = G = 1$

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$$h := 1 + (e + p) / \rho: \text{specific enthalpy, } e: \text{comoving internal energy density})$$

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Five equations with six unknowns ☺

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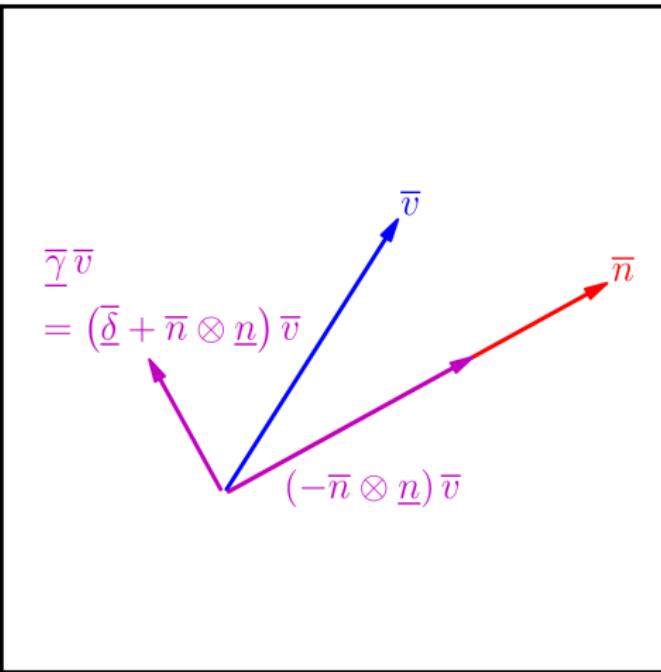
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Five equations with six unknowns ☺

Close with an equation of state: $p = p(e) := (\Gamma - 1) e$, $\Gamma = 4/3$

Valencia Decomposition



Extensible to higher-rank tensors!

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$$E := n_{\mu'} n_{\nu'} T^{\mu' \nu'}$$

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Math...

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$$\partial_t \mathbf{U} + \frac{1}{\sqrt{\gamma}} \partial_i [\alpha \sqrt{\gamma} \mathbf{F}^i (\mathbf{U})] = \mathbf{S} (\mathbf{U})$$

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$$\mathbf{F}^i (\mathbf{U}) = \begin{pmatrix} \rho W v^i \\ \rho h W^2 v^i v_j + p \delta^i_j \\ (\rho h W^2 - \rho W) v^i \end{pmatrix}$$

NR

$$\mathbf{F}^i (\mathbf{U}) = \begin{pmatrix} \rho v^i \\ \rho v^i v_j + p \delta^i_j \\ (\rho h_{\text{NR}} + \frac{1}{2} v^j v_j) v^i \end{pmatrix}$$

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$$\mathbf{S}(\mathbf{U}) = \begin{pmatrix} 0 \\ \frac{1}{2} \alpha P^{ik} \partial_j \gamma_{ik} - E \partial_j \alpha \\ -S^j \partial_j \alpha \end{pmatrix}$$

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$$\mathbf{S}(\mathbf{U}) = \begin{pmatrix} 0 \\ \frac{1}{2} P^{ik} \partial_j \bar{\gamma}_{ik} - \rho \partial_j \Phi \\ -S^j \partial_j \Phi \end{pmatrix}$$
$$\Phi(r) := -M/r$$

Steady-State Solutions (pre-shock)

$$\partial_t \vec{U}^0 + \frac{1}{\sqrt{\gamma}} \partial_r [\alpha \sqrt{\gamma} \vec{F}^r (\vec{U})] = \vec{S} (\vec{U})$$

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$$\frac{1}{2} v^2 + h_{\text{NR}} + \Phi = \mathcal{B}_{\text{NR}}$$

Steady-State Solutions (pre-shock)

$$\partial_t \vec{U}^0 + \frac{1}{\sqrt{\gamma}} \partial_r [\alpha \sqrt{\gamma} \vec{F}^r (\vec{U})] = \vec{S} (\vec{U})$$

GR

$$\alpha \psi^6 W \times 4\pi r^2 \rho v = -\dot{M}$$

$$\alpha h W = \mathcal{B}$$

$$p = K_1 \rho^\Gamma$$

NR

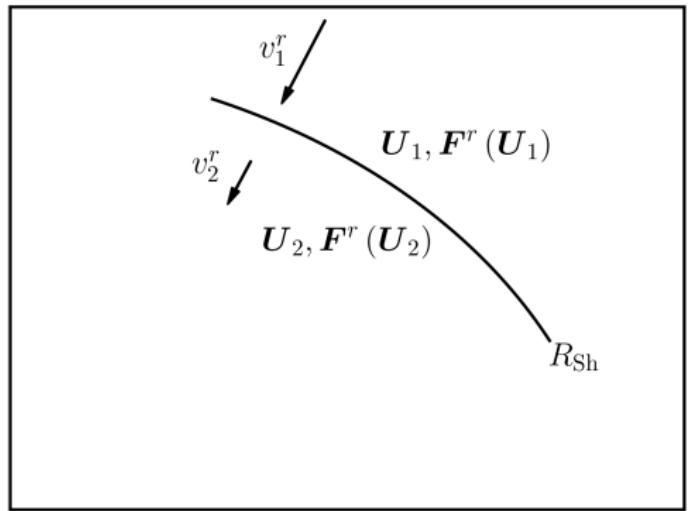
$$4\pi r^2 \rho v = -\dot{M}$$

$$\frac{1}{2} v^2 + h_{\text{NR}} + \Phi = \mathcal{B}_{\text{NR}}$$

$$p = K_1 \rho^\Gamma$$

Jump Conditions

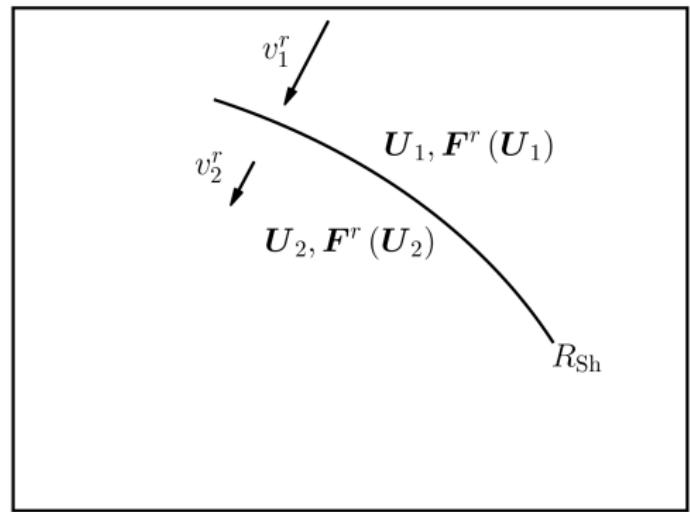
$$U_1 \neq U_2$$



Jump Conditions

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$$\mathbf{F}^r(U_1) = \mathbf{F}^r(U_2)$$

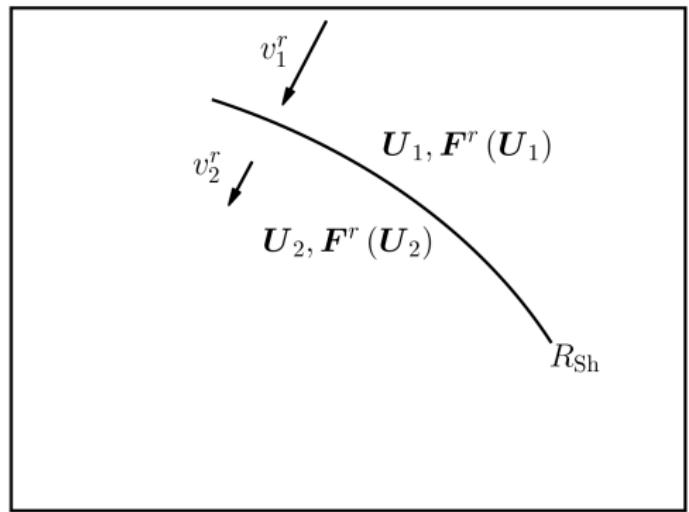


Jump Conditions

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Yields:



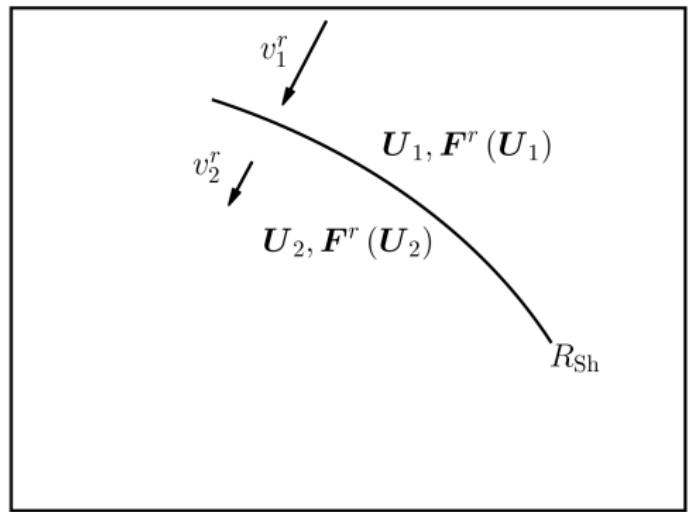
Jump Conditions

$$U_1 \neq U_2$$

$$\mathbf{F}^r(\mathbf{U}_1) = \mathbf{F}^r(\mathbf{U}_2)$$

Yields:

$$\rho_2 > \rho_1$$



Jump Conditions

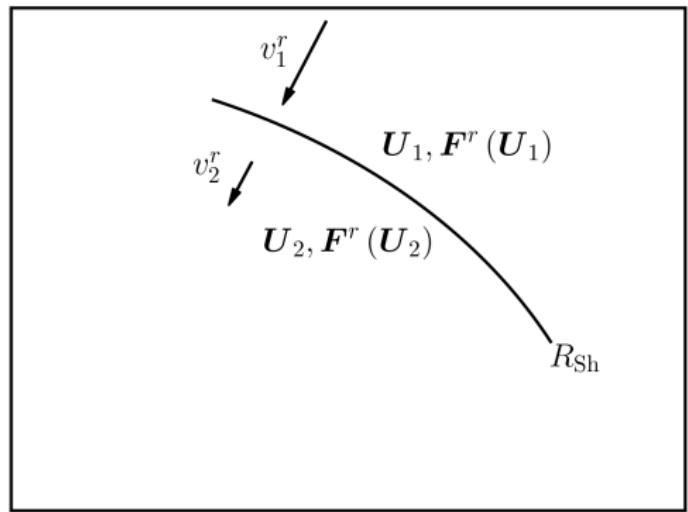
$$U_1 \neq U_2$$

$$\mathbf{F}^r(\mathbf{U}_1) = \mathbf{F}^r(\mathbf{U}_2)$$

Yields:

$$\rho_2 > \rho_1$$

$$e_2(p_2) > e_1(p_1)$$



Jump Conditions

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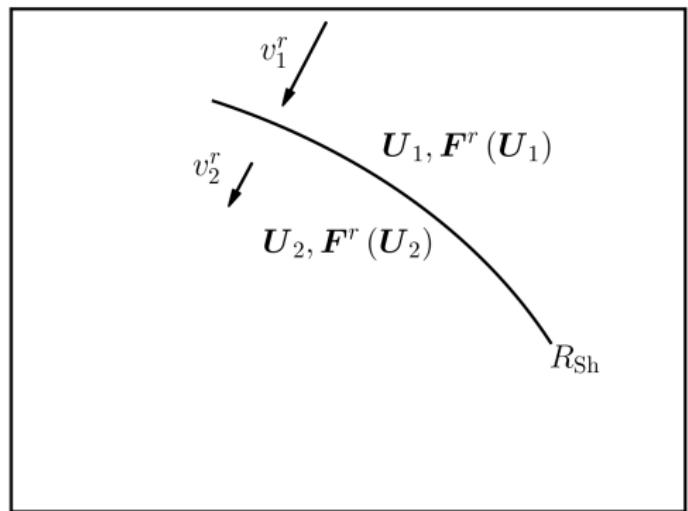
$$\mathbf{F}^r(U_1) = \mathbf{F}^r(U_2)$$

Yields:

$$\rho_2 > \rho_1$$

$$e_2(p_2) > e_1(p_1)$$

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Jump Conditions

$$\mathbf{U}_1 \neq \mathbf{U}_2$$

$$\mathbf{F}^r(\mathbf{U}_1) = \mathbf{F}^r(\mathbf{U}_2)$$

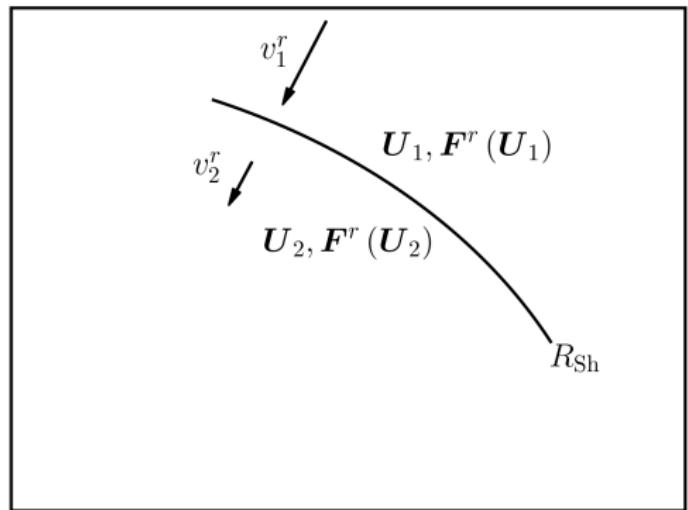
Yields:

$$\rho_2 > \rho_1$$

$$e_2(p_2) > e_1(p_1)$$

$$K_2 > K_1$$

$$|v_2^r| < |v_1^r|$$



Steady-State Solutions (post-shock)

$$\partial_t \vec{U}^0 + \frac{1}{\sqrt{\gamma}} \partial_r [\alpha \sqrt{\gamma} \vec{F}^r (\vec{U})] = \vec{S} (\vec{U})$$

GR

$$\alpha \psi^6 W \times 4\pi r^2 \rho v = -\dot{M}$$

$$\alpha h W = \mathcal{B}$$

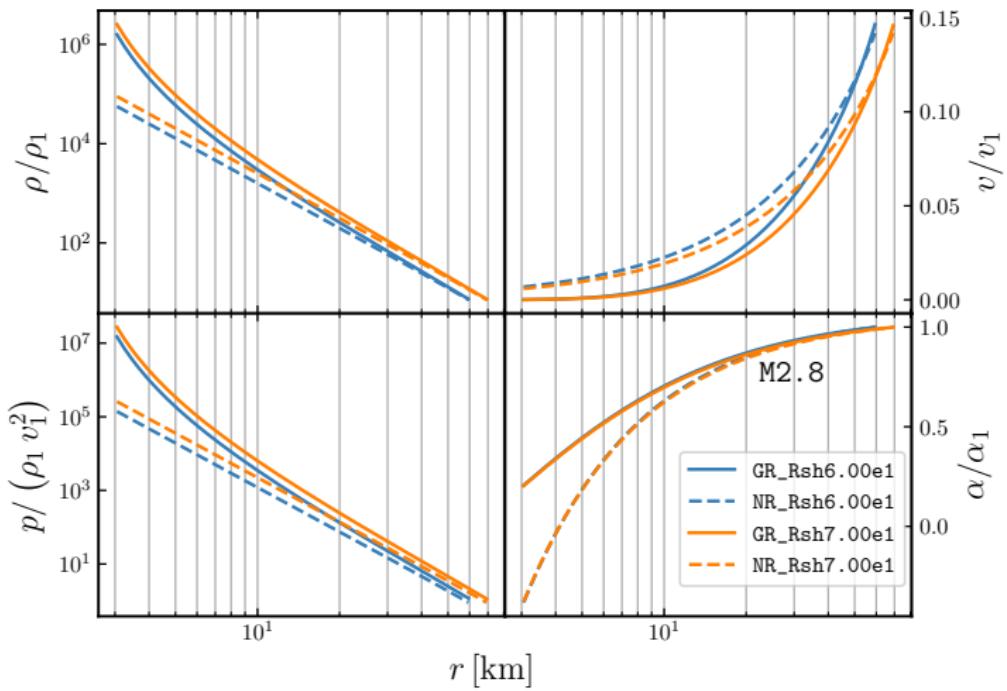
$$p = \textcolor{red}{K}_2 \rho^\Gamma$$

NR

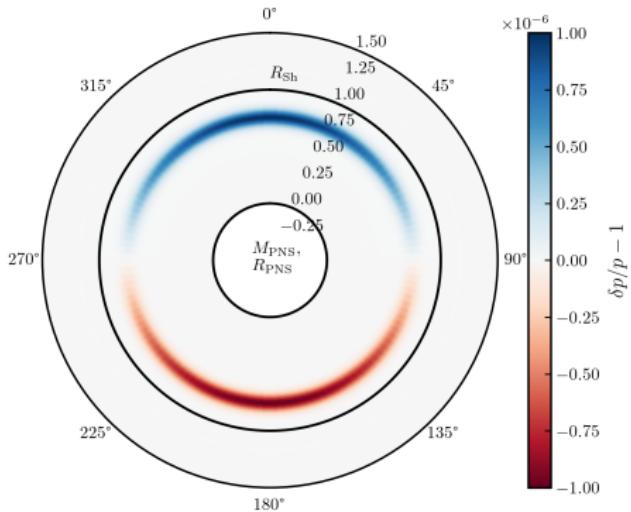
$$4\pi r^2 \rho v = -\dot{M}$$

$$\frac{1}{2} v^2 + h_{\text{NR}} + \Phi = \mathcal{B}_{\text{NR}}$$

$$p = \textcolor{red}{K}_2 \rho^\Gamma$$



$$\eta(r) := \frac{r - R_{\text{PNS}}}{R_{\text{Sh}} - R_{\text{PNS}}}$$

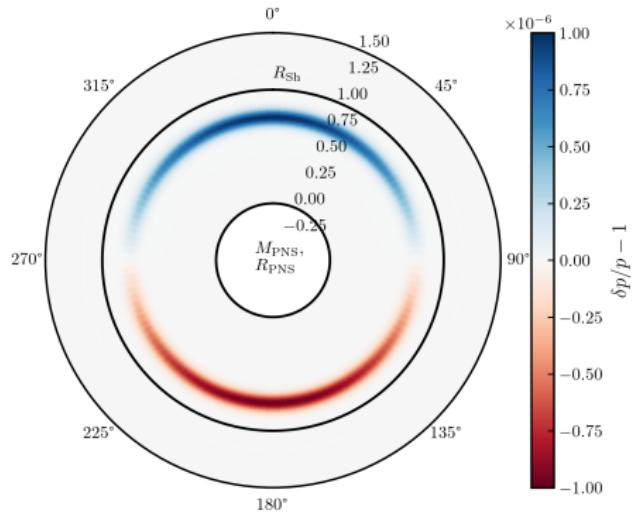


$$\eta(r) := \frac{r - R_{\text{PNS}}}{R_{\text{Sh}} - R_{\text{PNS}}}$$

$$\begin{aligned} \frac{\delta p(\eta, \theta)}{p(\eta_c)} &= \\ 10^{-6} \times \exp \left[\frac{-(\eta - \eta_c)^2}{2\sigma^2} \right] \cos \theta \end{aligned}$$

$$\eta_c = 0.75$$

$$\sigma = 0.05$$



Parameter Space

Model parameters:

M_{PNS} , R_{PNS} ,
 $R_{\text{Sh}}(t = 0)$, \dot{M} , K_1 ,
 $\mathcal{B}_{(\text{NR})}$, Γ

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$\xi =$
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2011), $R_{\text{Sh}}(t = 0)$

Table 1. Model Parameters

Model	$M_{\text{PNS}} [M_{\odot}]$	$R_{\text{PNS}} [\text{km}]$	$R_{\text{sh}} [\text{km}]$	ξ
M1.4_Rpns040_Rsh1.20e2	1.4	40	120	0.7
M1.4_Rpns040_Rsh1.50e2	1.4	40	150	0.7
M1.4_Rpns040_Rsh1.75e2	1.4	40	175	0.7
M2.8_Rpns020_Rsh6.00e1	2.8	20	60	2.8
M2.8_Rpns020_Rsh7.00e1	2.8	20	70	2.8

NOTE—Model parameters chosen for the 5 models. All models were run with both GR and NR. The first three rows correspond to the low-compactness models and the last two rows correspond to the high-compactness models.

$$A := \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} (v^\theta \sin \theta) \text{ (Scheck et al., 2008) } <\text{movie}>$$

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$$A(r, \theta, t) = \sum_{\ell'=0}^{\infty} G_{\ell'}(r, t) P_{\ell'}(\cos \theta)$$

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$$A(r, \theta, t) = \sum_{\ell'=0}^{\infty} G_{\ell'}(r, t) P_{\ell'}(\cos \theta)$$

$$\implies G_\ell(r, t) := \frac{1}{N_\ell} \int_0^\pi A(r, \theta, t) P_\ell(\cos \theta) \sin \theta d\theta$$

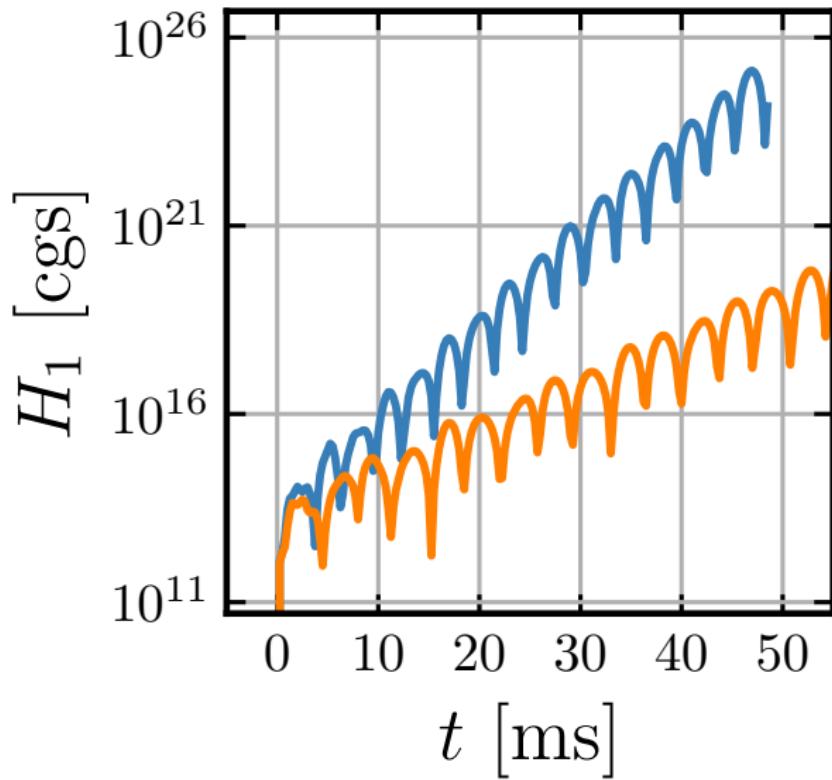
$$A := \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} (\sin \theta) \text{ (Scheck et al., 2008)} <\!\!\text{movie}\!>$$

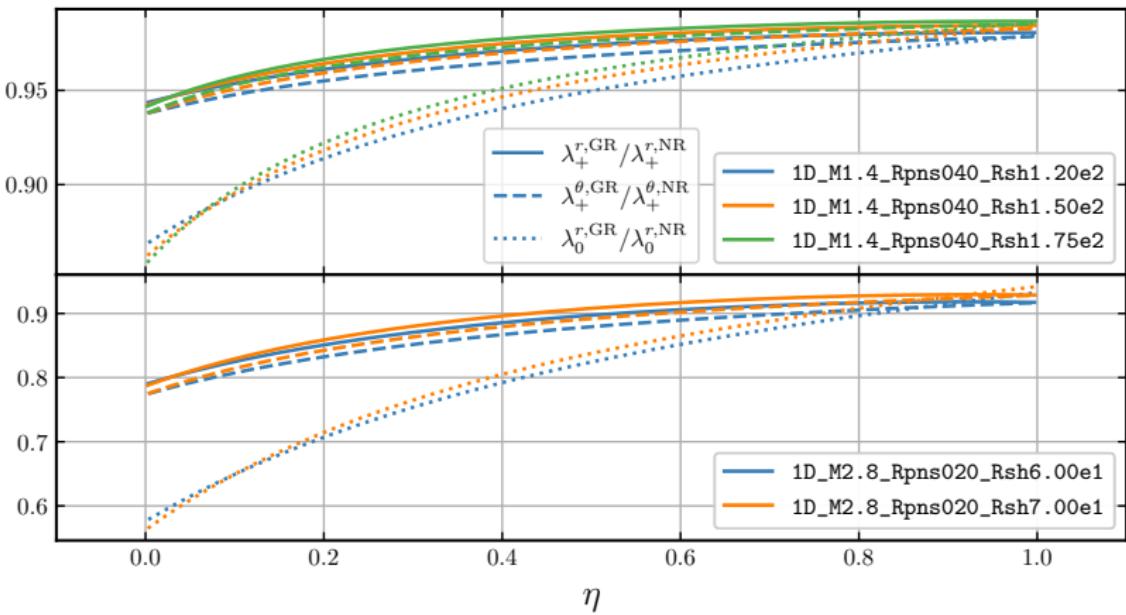
$$A(r, \theta, t) = \sum_{\ell'=0}^{\infty} G_{\ell'}(r, t) P_{\ell'}(\cos \theta)$$

$$\implies G_{\ell}(r, t) := \frac{1}{N_{\ell}} \int_0^{\pi} A(r, \theta, t) P_{\ell}(\cos \theta) \sin \theta d\theta$$

$$H_{\ell}(t) := 4\pi \int_{r_a}^{r_b} [G_{\ell}(r, t)]^2 [\psi(r)]^6 r^2 dr <\!\!\text{movie}\!>$$

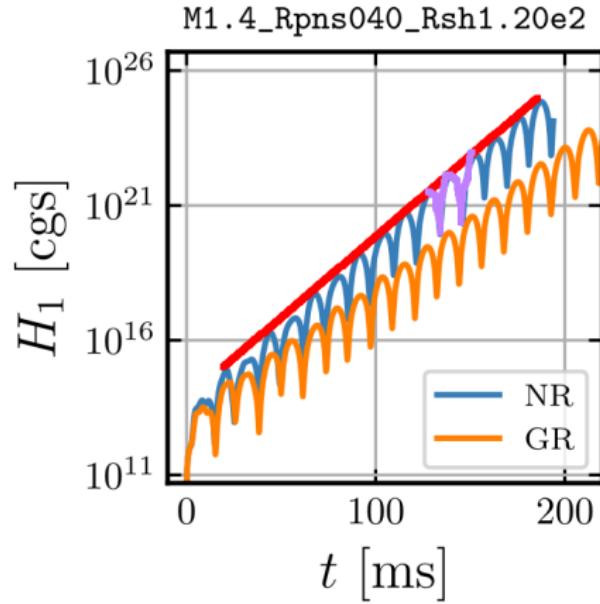
M2.8_Rpns020_Rsh6.00e1





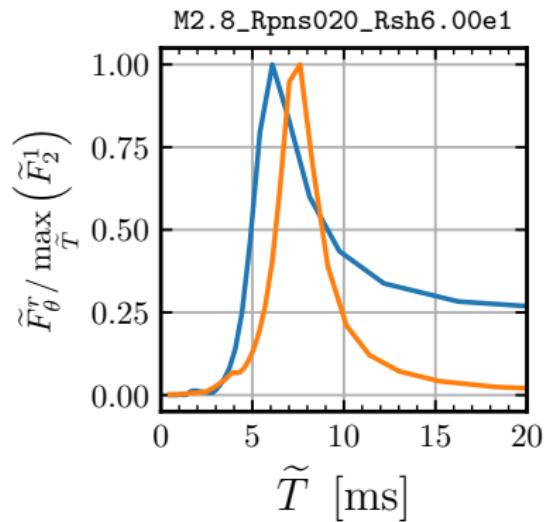
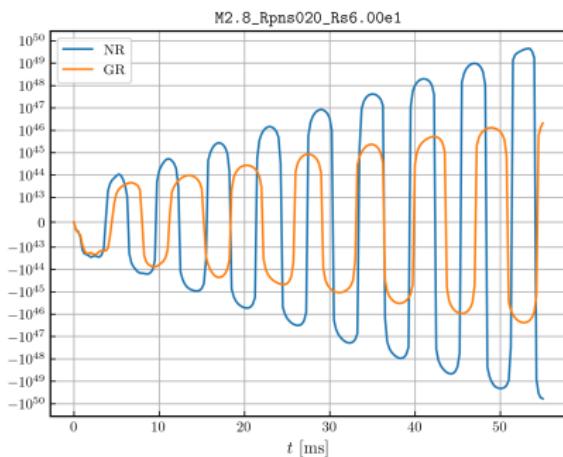
$$F(t) = F(0) e^{2\omega t} \sin^2 \left(\frac{2\pi t}{T} + \delta \right)$$

(Blondin and Mezzacappa, 2006)

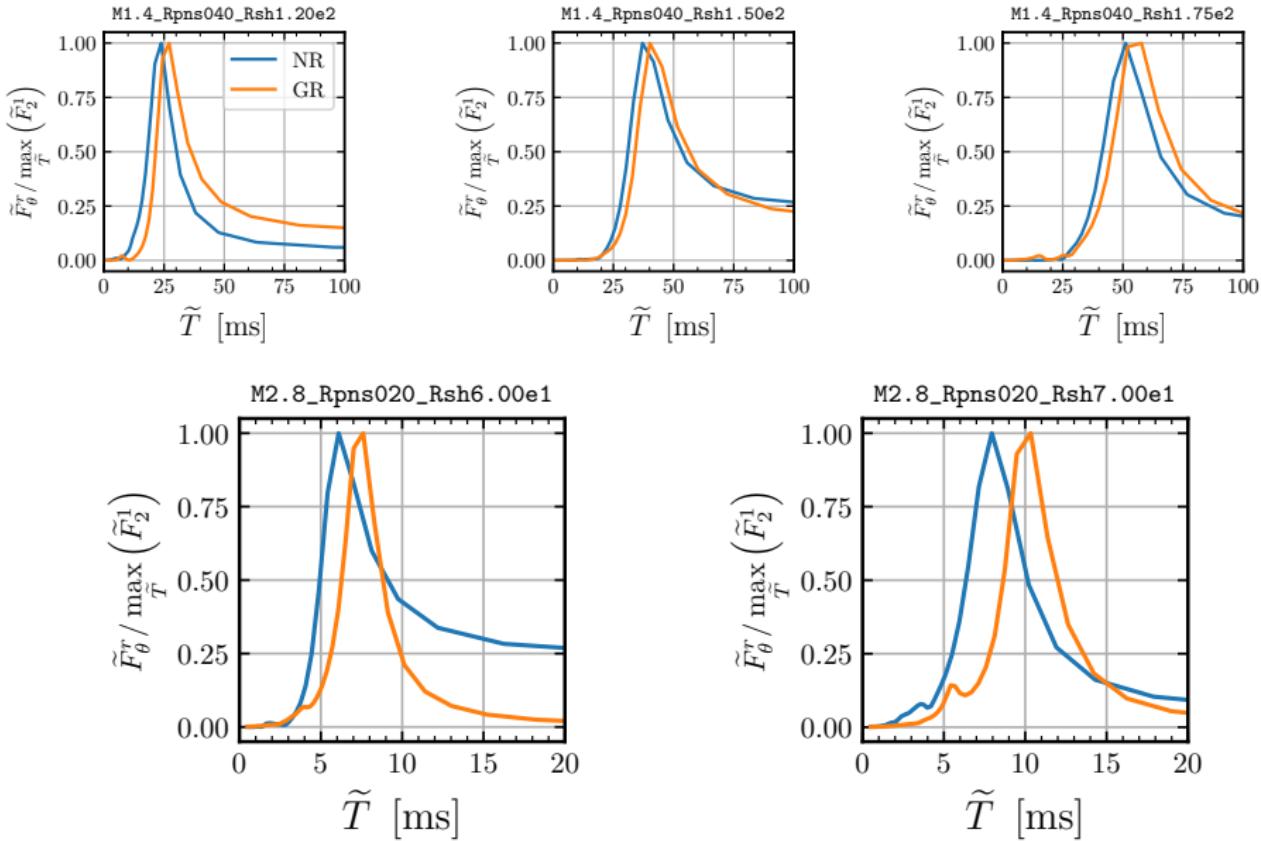


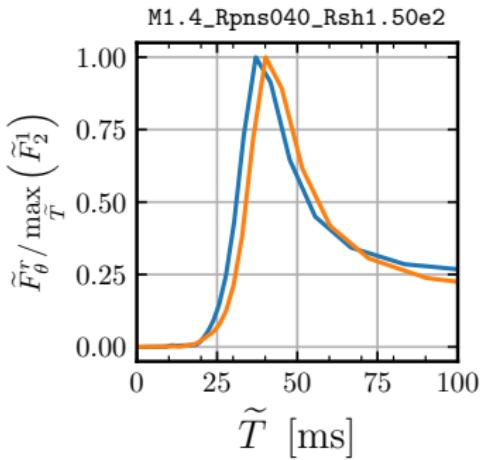
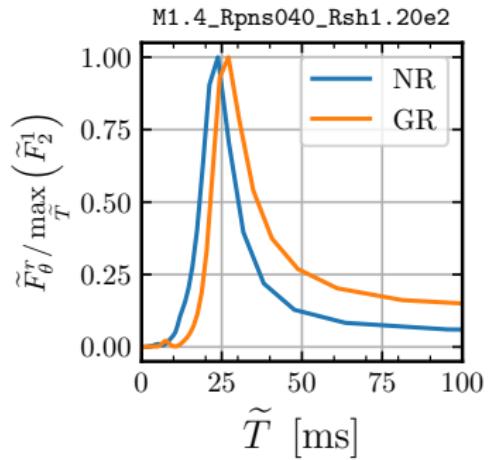
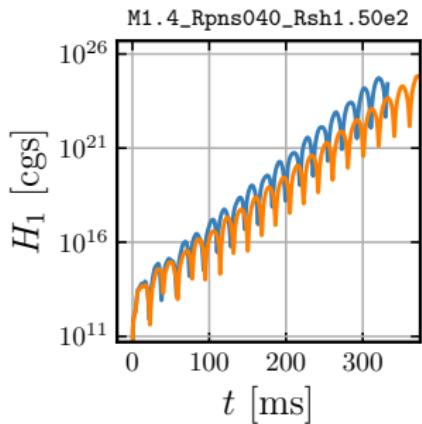
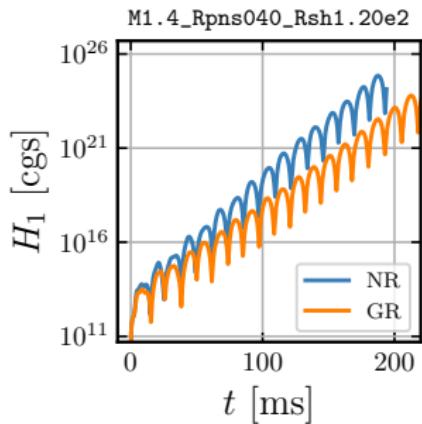
$$F_\theta^r := \alpha \psi^6 h W^2 \times \sqrt{\bar{\gamma}} \rho v^r v_\theta$$

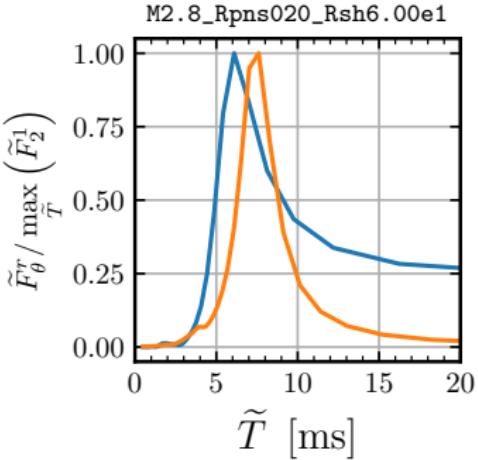
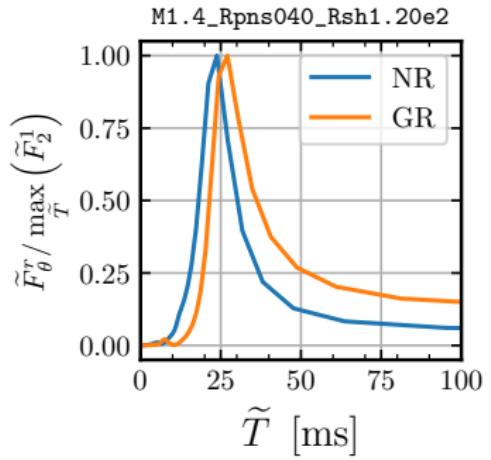
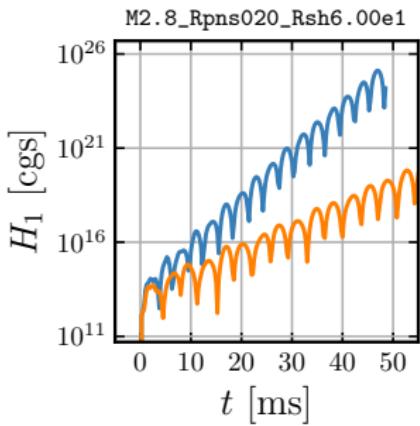
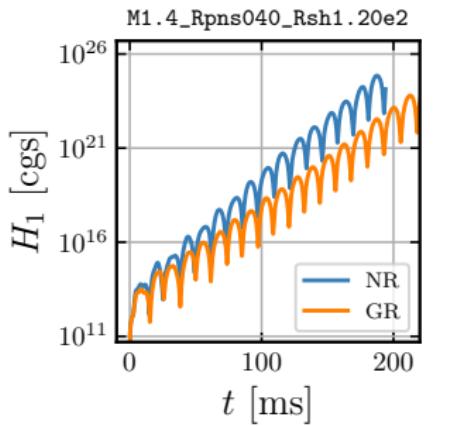
$$\tilde{F}_\theta^r = \text{FFT} \{ F_\theta^r \}$$

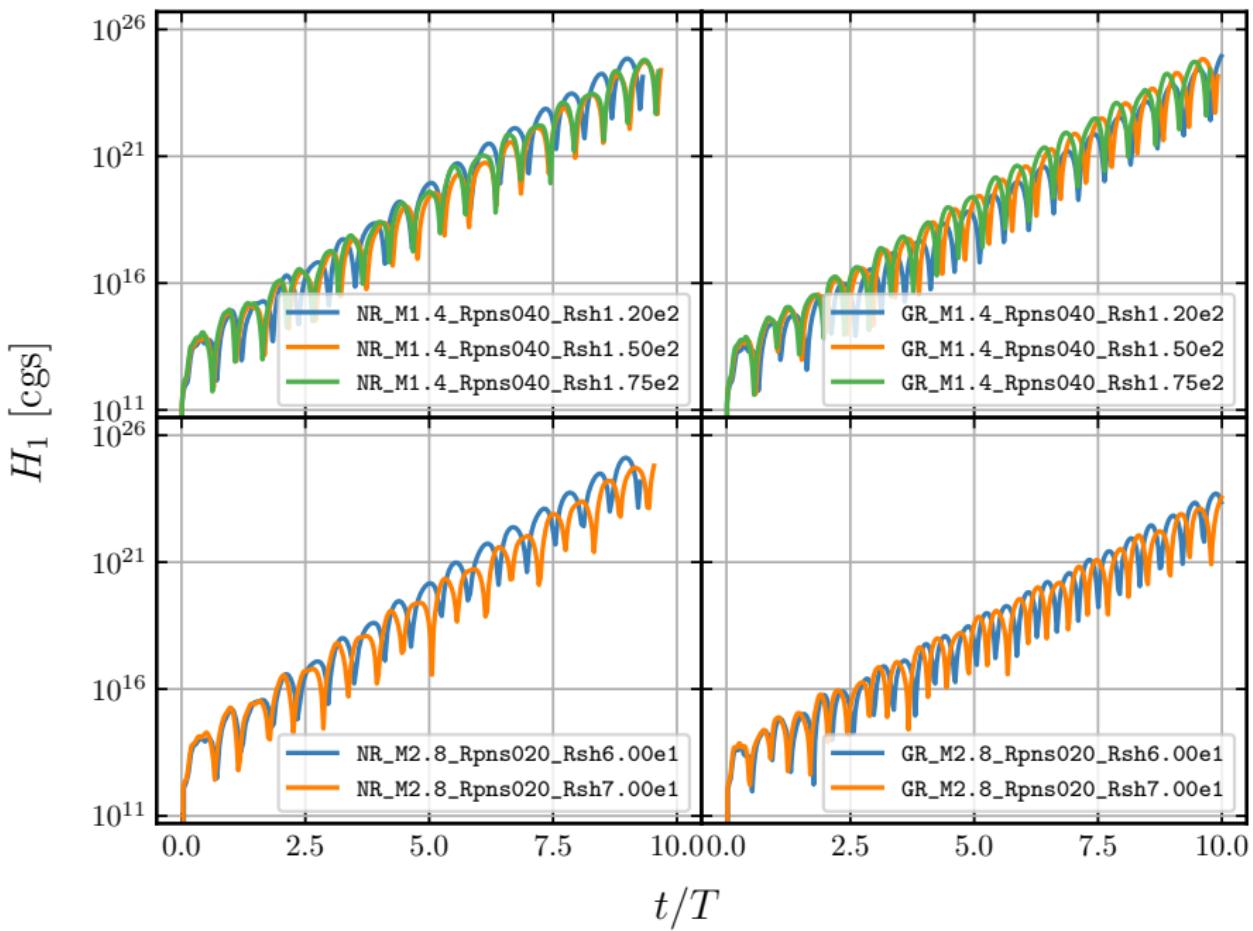


T defined as the unique \tilde{T} such that $\tilde{F}_\theta^r (\tilde{T}) = 1$

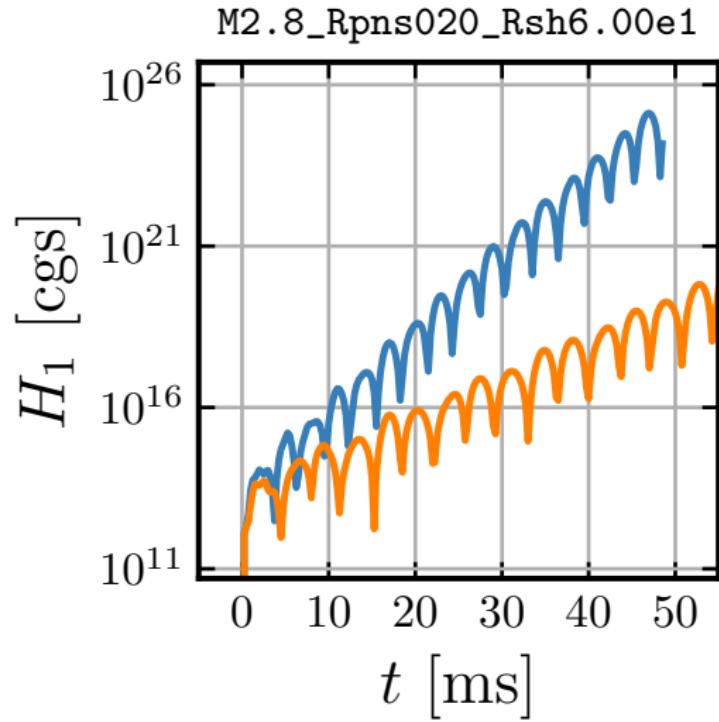




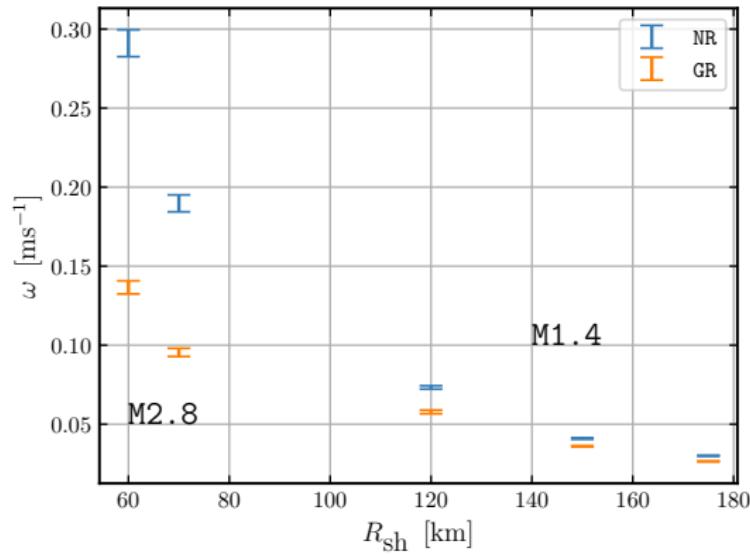




Conclusions



Conclusions



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Summary

- Extended study of Blondin and Mezzacappa (2006) to include GR
- Showed that GR leads to longer SASI oscillation period than NR
- Showed that GR leads to smaller SASI growth rate than NR
- Found that growth rate is such that ωT is roughly constant for some parameter sets: implications for growth rate mechanism
- Future Work
 - Extend study to 3D
 - Include GR monopole (Marek et al., 2006)